

# Optimal free models of many-body interacting theories

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## Motivation

- Interacting particles give rise to many fascinating phases of matter which remain theoretically challenging. Free systems are in contrast much better understood.
- Interacting systems are often described in terms of free and integrable models – as in mean field theory and perturbation theory.
- We quantify the relevance of interactions based on the particular structure of entanglement found in free systems.
- We use this measure to identify *optimal* free descriptions of interacting systems.

## Interaction Distance

We introduce the *interaction distance* between an interacting state  $\rho$  and manifold of free states  $\mathcal{F}$

$$T_{\mathcal{F}}(\rho) = \min_{\sigma \in \mathcal{F}} T(\rho, \sigma), \quad (1)$$

where  $T(\rho, \sigma) = \frac{1}{2} \text{tr} \sqrt{(\rho - \sigma)^2}$  is the trace distance. It measures the distinguishability of  $\rho$  from ground states of free theories.

The manifold  $\mathcal{F}$  contains all unitary orbits of Gaussian states, or equivalently Gaussian states in any basis of mode operators  $\{c\}$ . In particular,  $\sigma$  can be free in terms of quasiparticles with different exchange statistics than those of  $\rho$ .

We have freedom to consider general unitary orbits because each  $\sigma$  can be expressed as

$$\sigma = \exp \left\{ z + \sum_j m_j c_j^\dagger c_j \right\} \quad (2)$$

with some  $\{c\}$  bosonic or fermionic mode operators. The action of a general unitary operator  $U$  on  $\sigma$  simply effects a canonical transformation  $c \mapsto UcU^\dagger$  leaving the system free even if it is not a Gaussian map.

It has been shown that the trace distance is minimised within a unitary orbit when  $\sigma$  and  $\rho$  are simultaneously diagonal and their eigenvalues are both in rank order. Only minimisation with respect to  $\{\epsilon\}$  is needed in order to compute interaction distance which makes this procedure efficient.

## Free entanglement spectra

The reduced density matrix  $\rho$  for a ground state  $|\psi\rangle$  is  $\rho = \text{tr}_B |\psi\rangle \langle \psi|$  found by tracing out part of the system. The entanglement Hamiltonian  $H_E = -\ln \rho$  has eigenvalues  $\{\xi\}$ , known as the *entanglement spectrum*.

The entanglement spectrum of a free state is generated combinatorially from the *single-body* entanglement spectrum  $\{\epsilon\}$ ,

$$\xi_n = \left\{ \xi_0 + \sum_{i=1}^N n_i \epsilon_i \right\}_{n_i=0, \dots, m} \quad (3)$$

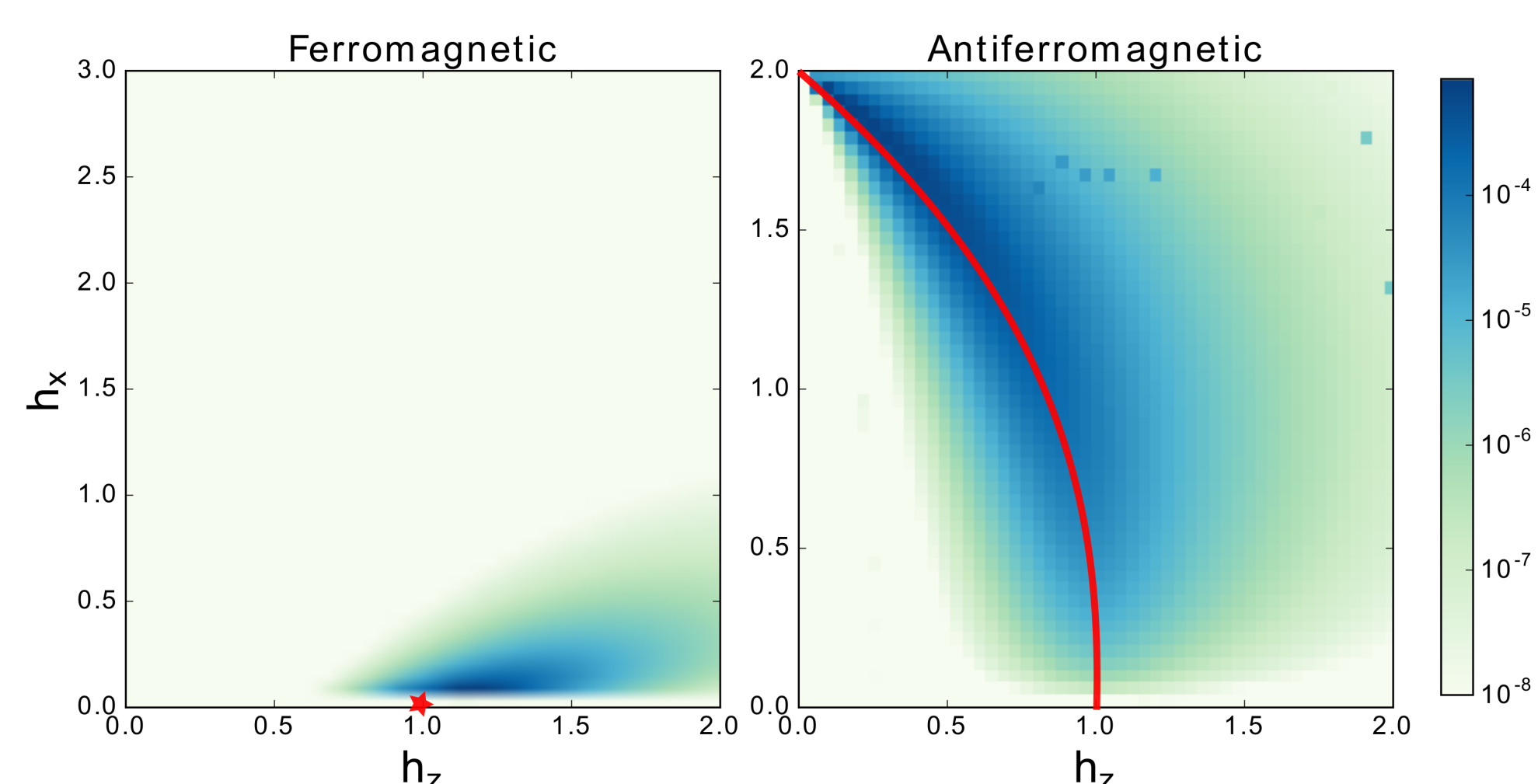
with  $m = \infty$  for bosons,  $m < \infty$  for soft-core bosons, and  $m = 1$  for fermions, and  $\xi_0$  a normalisation constant. This is because the partial trace is a Gaussian map.

## Case study: Ising model in a magnetic field

As an example we use the 1D quantum ferromagnetic (FM) and antiferromagnetic (AFM) Ising models in transverse ( $h_x$ ) and ( $h_z$ ) longitudinal fields

$$H_{\pm} = - \sum_j (\pm \sigma_j^x \sigma_{j+1}^x + h_z \sigma_j^z + h_x \sigma_j^x). \quad (4)$$

In both models the  $h_x = 0$  transverse Ising line is free and satisfies  $T_{\mathcal{F}} = 0$ . This line contains a critical point at  $h_x = 0, h_z = 1$  which is isolated in the FM model and connected by a critical line to the classical critical point  $h_z = 0, h_x = 2$  in the AFM model (sketched).



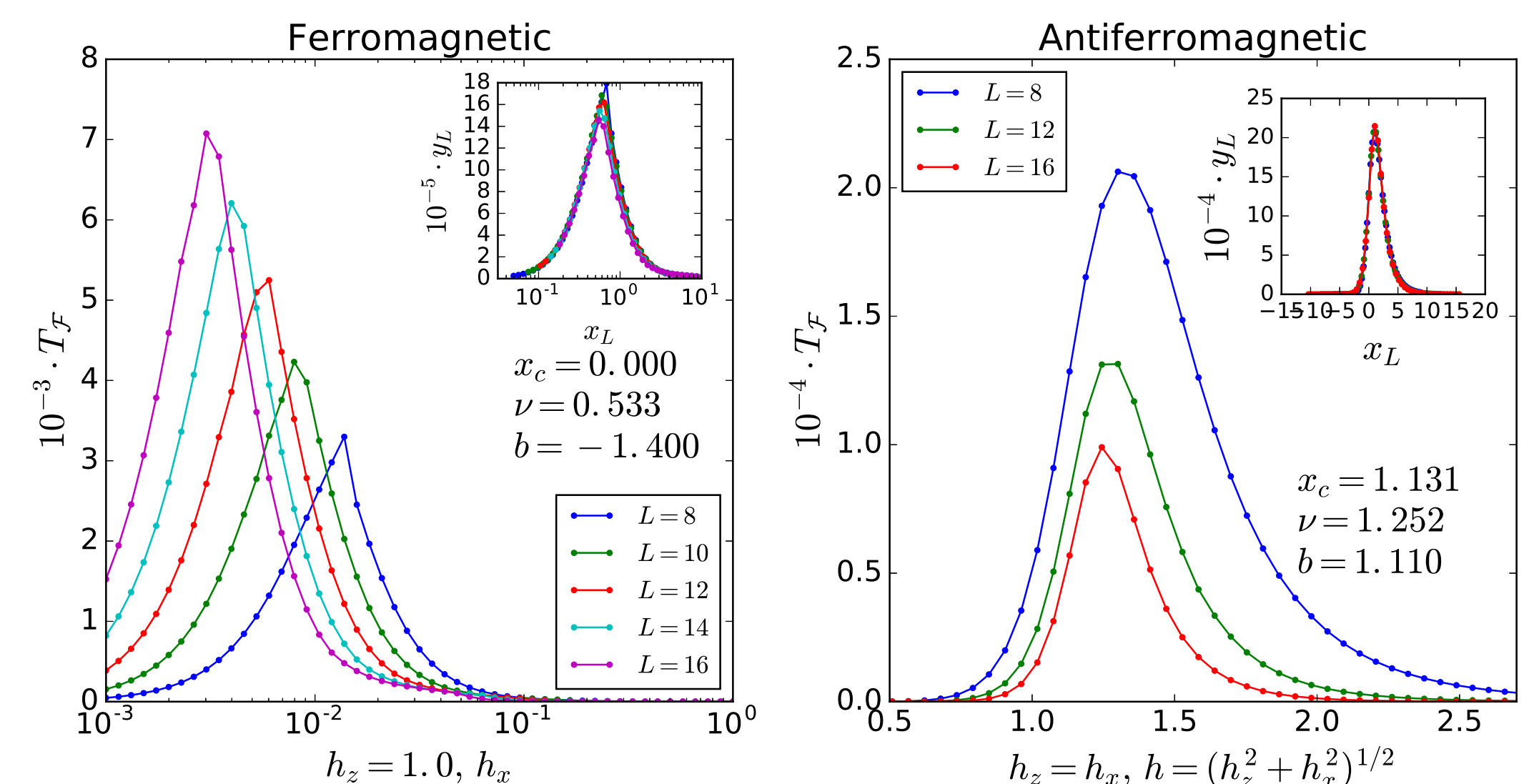
- The figures are for periodic boundary conditions and 16 spins.
- $T_{\mathcal{F}}$  is significant only close to criticality and decreasing toward zero approaching the stable fixed points of the renormalisation flow.
- As the system size is increased the region over which  $T_{\mathcal{F}}$  is significant shrinks around criticality.

## Scaling around critical points

We examine the scaling around criticality with finite size scaling (FSS) analysis according to

$$T_{\mathcal{F}}(\lambda; L) \sim L^{-b} \tilde{f} \left( (h - h_c) L^{1/\nu} \right) \quad (5)$$

The exponent  $\nu$  is the correlation length critical exponent. Sign of exponent  $b$  relates to the relevance or irrelevance of corresponding operator in the critical theory.



- From FSS analysis of the energy gap  $\Delta E$  we find  $\nu \approx 1.052$  for the AFM model which is within 20% of the value obtained from  $T_{\mathcal{F}}$ .
- We know analytically  $\nu$  for the FM model where it is  $\nu = 8/15$  which is in good agreement.

## Convergence

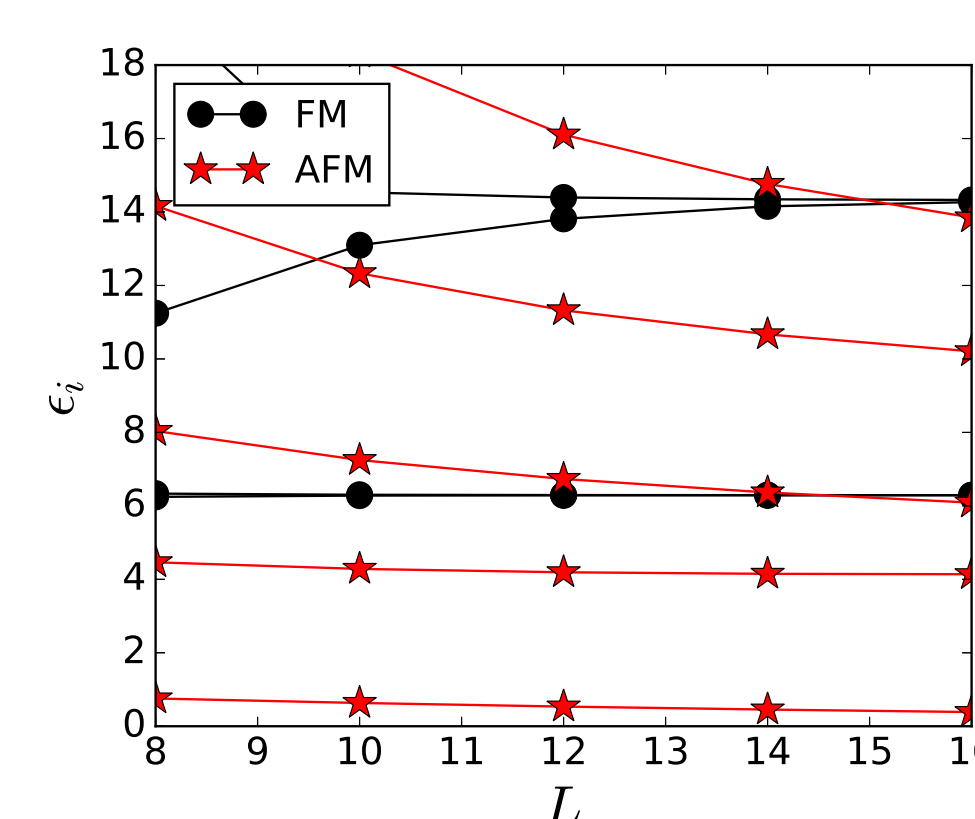
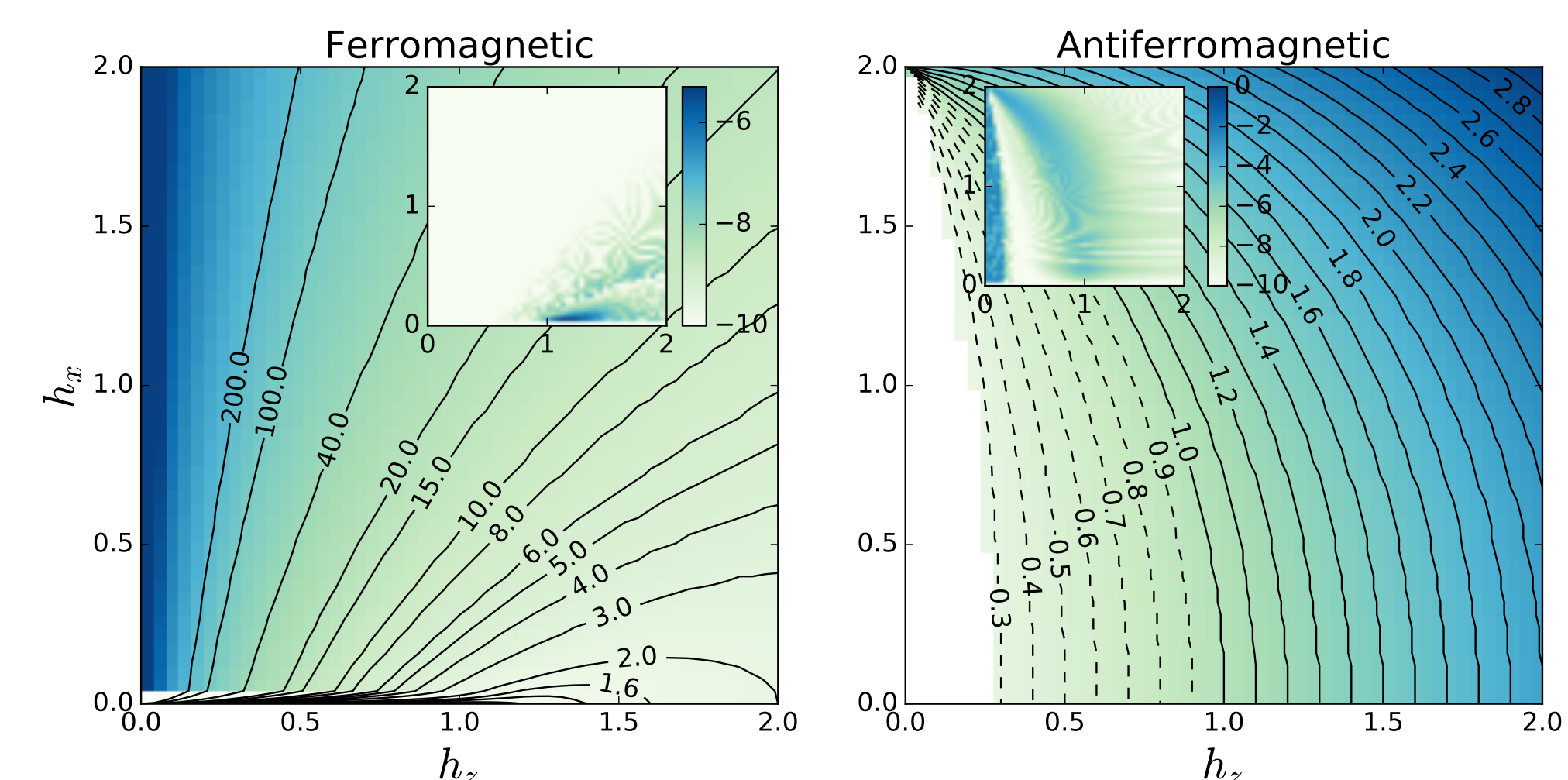


Figure shows the convergence the single particle entanglement energies  $\epsilon$  as a function of system size  $L$  to their asymptotic values in the thermodynamic limit for the (blue) FM and (red) AFM models at  $(h_x = 0.16, h_z = 0.88)$ .

## Matching the free model

We minimise the trace distance between the optimal free model and the free Ising models identifying the free theory ( $h_z^{\mathcal{F}}$ ) with equivalent ground state correlations to the interacting theory ( $h_z, h_x$ ).



- Contours shows matching of the optimal free model to the free Ising models in the thermodynamic limit.
- The Insets give the trace distance (log scale) between the optimal free model to the Ising free model.
- Unresolved degeneracy in the energy spectrum close to the classical axis  $h_z = 0$  causes artefacts which have been removed.

## Outlook

- Could be used to find a free parent Hamiltonian.
- The scaling exponents around criticality could be explained using conformal field theory.
- See what happens for non-renormalisable models.

## References

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